# A new set of generators and a physical interpretation for the SU(3) finite subgroup D(9, 1, 1; 2, 1, 1)

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#### Abstract

After 100 years of effort, the classification of all the finite subgroups of SU(3) is yet incomplete. The most recently updated list can be found in [3], where the structure of the series (C) and (D) of SU(3)-subgroups is studied. We provide a minimal set of generators for one of these groups which has order 162. These generators appear up to phase as the image of an irreducible unitary braid group representation issued from the Jones-Kauffman version of SU(2) Chern-Simons theory at level 4. In light of these new generators, we study the structure of the group in detail and recover the fact that it is isomorphic to the semidirect product  $\mathbb{Z}_9 \times \mathbb{Z}_3 \rtimes S_3$  with respect to conjugation.

## **1** Introduction and main result

### **1.1** Definition of the group

Over the past century, there has been interest in classifying and studying the structure of all the finite subgroups of SU(3) because these subgroups appear in physics. A recent paper by Patrick Otto Ludl [3] gives in its introduction a complete chronological review on the matter which we skip in the present paper. As of today, the classification of all the finite subgroups of SU(3) remains incomplete. The original classification attempt dates from 1916 in a work of Blichfeldt [4]. Two series of finite subgroups of SU(3)named (C) and (D) are then defined but the complete structure of some of these groups was not studied until 2011 in [3]. Our paper focuses on one particular group of type D of order 162, namely D(9, 1, 1; 2, 1, 1). The series (C) and (D) respectively contain subseries  $\Delta(3n^2)$  and  $\Delta(6n^2)$ . For a definition of these series, see for instance [2]. It was shown in 2009 by P.O. Ludl in his thesis [2] that any SU(3) subgroup of type (C) can be interpreted as a three-dimensional irreducible representation of  $\Delta(3n^2)$ . As shown by the same author, the group D(9, 1, 1; 2, 1, 1) provides a counter-example that the same kind of result does not hold for the SU(3) subgroups of type (D), that is there exists at least one SU(3) subgroup of type D that cannot be interpreted as a three-dimensional irreducible representation of  $\Delta(6n^2)$ . We recall below the definitions of the series (C) and (D) as in [3]. The series C(n, a, b) for n a positive integer and a and b integers with  $0 \le a, b \le n - 1$ , is defined as the group generated by the permutation matrix E corresponding to the cycle (1, 3, 2) of the symmetric group  $S_3$  and the diagonal matrix F(n, a, b) with n-th roots of unity on the diagonal and determinant 1.

$$C(n, a, b) = \langle E, F(n, a, b) \rangle$$
, where

$$E = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} \text{ and } F(n, a, b) = \begin{pmatrix} e^{\frac{2ia\pi}{n}} & 0 & 0 \\ 0 & e^{\frac{2ib\pi}{n}} & 0 \\ 0 & 0 & e^{\frac{2i(-a-b)\pi}{n}} \end{pmatrix}$$

In order to get the SU(3) subgroup D(n, a, b; d, r, s), with d a positive integer and r and s two integers with  $0 \le r, s \le d - 1$ , add the extra generator

$$\begin{pmatrix} e^{\frac{2ir\pi}{d}} & 0 & 0\\ 0 & 0 & e^{\frac{2is\pi}{d}}\\ 0 & -e^{\frac{2i(-r-s)\pi}{d}} & 0 \end{pmatrix}$$

## 1.2 The new generators

The interest of our paper is double. We are able to provide only two generators for D(9, 1, 1; 2, 1, 1) instead of three. Moreover, our two generators are issued from an irreducible unitary braid group representation  $B_4 \rightarrow U(3)$ obtained by braiding four anyons of topological charge 2 on a fusion tree of total topological charge 0 in the Jones-Kauffman version of SU(2) level 4 Chern-Simons theory [1]. For some explanations and some terminology associated with this theory, we refer the reader to the excellent exposition in [5]. We recover the structure of the group as a semidirect product of the normal abelian group  $\mathbb{Z}_9 \times \mathbb{Z}_3$  and the symmetric group on three letters  $S_3$ , where the action is given by conjugation. Our results are summarized below.

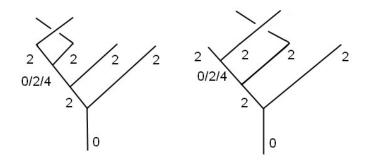
**Theorem 1.** Let  $t = \frac{\sqrt{2}}{2}e^{\frac{2i\pi}{3}}$  and let  $\mathcal{G}$  be the subgroup of SU(3) generated by the

matrices  $G_1$  and  $G_2$ , defined as follow.

$$G_{1} = e^{\frac{i\pi}{9}} \begin{pmatrix} 2\bar{t}^{2} & & \\ & 2t^{2} & \\ & & -2\bar{t}^{2} \end{pmatrix} \quad G_{2} = e^{\frac{i\pi}{9}} \begin{pmatrix} t^{2} & t & -t^{2} \\ t & 0 & t \\ -t^{2} & t & t^{2} \end{pmatrix}$$

The matrices  $G_1$  and  $G_2$  have the same eigenvalues and so  $G_2$  is unitarily similar  $\int \frac{i7\pi}{2}$ 

to 
$$\begin{pmatrix} e^{\overline{9}} & \\ & -e^{\frac{i4\pi}{9}} \\ & & e^{\frac{-2i\pi}{9}} \end{pmatrix}$$
. Both matrices thus have order 18.



Up to the phase  $e^{\frac{i\pi}{9}}$ , the matrices  $G_1$  and  $G_2$  are obtained by respectively braiding the anyons 1 and 2 on one hand and 2 and 3 on the other hand, like on the above two trees, where these are subject to adequate unitary normalizations. The matrices  $G_1$  and  $G_2$  satisfy the braid relation

$$G_1 G_2 G_1 = G_2 G_1 G_2$$

and their respective squares commute

$$G_1^2 G_2^2 = G_2^2 G_1^2$$

 $\mathcal{G} = \mathcal{N} \stackrel{\phi}{\rtimes} \mathcal{H}$ , where the normal subgroup  $\mathcal{N}$  of  $\mathcal{G}$  and the subgroup  $\mathcal{H}$  of  $\mathcal{G}$  are respectively defined by

$$\mathcal{N} = < G_1 \, G_2^2 \, G_1^{-1}, \, G_1 \, G_2^{-2} \, G_1 > \simeq \mathbb{Z}_9 \times \mathbb{Z}_3$$

$$\mathcal{H} = \{I, \begin{pmatrix} -1 & & \\ & -1 & \\ & & 1 \end{pmatrix}, G_1 G_2 G_1, \begin{pmatrix} -1 & & \\ & -1 & \\ & & 1 \end{pmatrix} G_1 G_2 G_1 \begin{pmatrix} -1 & & \\ & & -1 & \\ & & 1 \end{pmatrix}, G_1 G_2 G_1 \begin{pmatrix} -1 & & \\ & & 1 \end{pmatrix}, G_1 G_2 G_1 \} \simeq S_3,$$

and where for any matrix H of  $\mathcal{H}$ , the map  $\phi(H)$  is the automorphism of  $\mathcal{N}$  which is the conjugation by H. Thus  $\mathcal{G}$  is isomorphic to D(9, 1, 1; 2, 1, 1).

The non-diagonal matrices defining  $\mathcal{H}$  are, in the same order as given above,

$$\begin{pmatrix} -\frac{1}{2} & -\frac{1}{\sqrt{2}} & -\frac{1}{2} \\ -\frac{1}{\sqrt{2}} & 0 & \frac{1}{\sqrt{2}} \\ -\frac{1}{2} & \frac{1}{\sqrt{2}} & -\frac{1}{2} \end{pmatrix}, \begin{pmatrix} -\frac{1}{2} & -\frac{1}{\sqrt{2}} & \frac{1}{2} \\ -\frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{\sqrt{2}} & -\frac{1}{2} \end{pmatrix}, \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} & -\frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & \frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{\sqrt{2}} & -\frac{1}{2} \end{pmatrix}, \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{\sqrt{2}} & -\frac{1}{2} \end{pmatrix}, \begin{pmatrix} \frac{1}{2} & \frac{1}{\sqrt{2}} & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ -\frac{1}{2} & \frac{1}{\sqrt{2}} & -\frac{1}{2} \end{pmatrix}$$

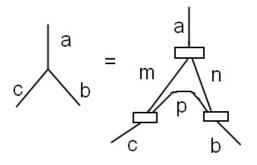
The paper is organized as follows. We define the Hilbert space of the unitary braid group representation whose image is the group that we are studying here and recall some facts in Temperley-Lieb recoupling theory. We then introduce the two generators of the abelian subgroup  $\mathcal{N}$  and show that this subgroup is a normal subgroup of  $\mathcal{G} = \langle G_1, G_2 \rangle$ . We next prove that  $\mathcal{G}$  is the semidirect product announced in Theorem 1.

# 2 The unitary braid group representation

In what follows,  $d = -A^2 - A^{-2}$  is the loop variable in the Temperley-Lieb algebra and the Kauffman variable A is set to the value  $A = i e^{-\frac{2i\pi}{4r}}$ , where k = r - 2 is the level of the theory. Here we work at level 4 and so  $A = i e^{-\frac{i\pi}{12}}$ . The set  $L = \{0, 1, 2, 3, 4\}$  is the label set of the different particule types. The particules obey fusion rules such that the labels *a*, *b* and *c* of a trivalent vertex must satisfy

$$\left\{\begin{array}{l}a+b+c \text{ is even}\\c\leq a+b,\ b\leq a+c,\ a\leq b+c\\a+b+c\leq 2k\end{array}\right.$$

Such a triple (a, b, c) is said to be admissible. By definition,



with

$$m = \frac{a+c-b}{2}, \ n = \frac{a+b-c}{2}, \ p = \frac{b+c-a}{2}$$

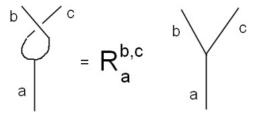
and where the boxes to the right of the figure represent Jones-Wenzl projectors. We define  $\Delta_n$  as the bracket evaluation of the closure of the (n-1)-th projector and so  $\Delta_1 = d$ . The classical recursion formula for the Jones-Wenzl projectors shows that  $\Delta_n$  is a Chebyschev polynomial of the second kind, whence

$$\Delta_n = (-1)^n \left[ n+1 \right]$$

In this formula, [n] denotes the quantum integer taken at  $q = A^2$ .

$$[n] = \frac{\mathsf{A}^{2n} - \mathsf{A}^{-2n}}{\mathsf{A}^2 - \mathsf{A}^{-2}}$$

We have in particular  $\Delta_0 = \Delta_4 = 1$  and  $\Delta_2 = 2$ . We recall below the *R*-move followed by the *F*-move. The theory appears in [1].



with

$$R_a^{b,c} = (-1)^{\frac{b+c-a}{2}} \mathsf{A}^{\frac{b(b+c)+c(c+2)-a(a+2)}{2}}$$

In our forthcoming computations, a few useful values are the following.  $\overline{R_0^{2,2}} = e^{\frac{2i\pi}{3}}, \overline{R_2^{2,2}} = -e^{\frac{i\pi}{3}}, \overline{R_4^{2,2}} = e^{-\frac{i\pi}{3}}$ , where the bar denotes the complex conjugate.

,

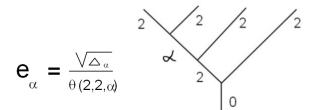
where k runs over all the labels such that both triples (b, c, k) and (a, d, k) are admissible, and where the braces denote the 6j-symbol

$$\left\{ \begin{array}{cc} a & b & k \\ c & d & i \end{array} \right\} = \frac{T \left[ \begin{array}{cc} a & b & k \\ c & d & i \end{array} \right] \Delta_i}{\theta(a, d, i)\theta(c, b, k)}$$
(1)

For the evaluations of the tetrahedral net T and the theta net  $\theta$  in terms of the quantum integers, and other formulas from recoupling theory, we refer the reader to the summary in Chapter 9 of [1]. We will frequently use that

$$T\begin{bmatrix} 2 & 2 & j \\ 2 & 2 & i \end{bmatrix} = T\begin{bmatrix} 2 & 2 & i \\ 2 & 2 & j \end{bmatrix} = \begin{bmatrix} 2 & \text{if } (i,j) = (0,0), \\ \frac{2}{\sqrt{3}} & \text{if } (i,j) = (2,0), \\ 0 & \text{if } (i,j) = (2,2), \\ 1 & \text{if } (i,j) = (4,0), \\ -\frac{1}{\sqrt{3}} & \text{if } (i,j) = (4,2), \\ \frac{1}{2} & \text{if } (i,j) = (4,4) \end{bmatrix}$$

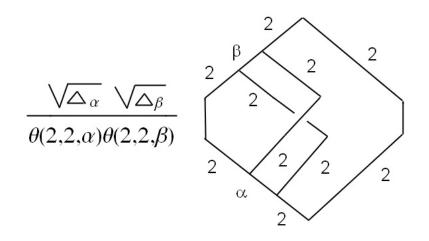
We consider the 3-dimensional vector space  $V_{4,2,0}$  spanned over  $\mathbb{C}$  by the vectors  $e_{\alpha}$  with  $\alpha \in \{0, 2, 4\}$ .



We define an inner product over  $V_{4,2,0}$  so that the vectors  $e_{\alpha}$ ,  $\alpha = 0, 2, 4$ , form an orthonormal basis of  $V_{4,2,0}$ . The inner product between two such vectors  $e_{\beta}$  and  $e_{\gamma}$  is obtained by stacking the mirror image of the tree in  $e_{\beta}$  on top of the tree in  $e_{\gamma}$  and by resolving the crossings with the Kauffman bracket. We now consider the braid group  $B_4$  on four strands and we denote by  $g_1$ ,  $g_2$  and  $g_3$  its three generators. We see that the actions by  $g_1$ and  $g_3$  on  $e_{\alpha}$  can be resolved using an *R*-move and we get in both cases the same matrix

$$\begin{pmatrix} \overline{R_0^{2,2}} & & \\ & \overline{R_2^{2,2}} & \\ & & \overline{R_4^{2,2}} \end{pmatrix}$$

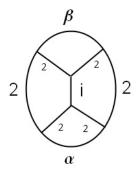
This matrix is the unitary diagonal matrix  $G_1$  of Theorem 1 before the phase  $e^{\frac{i\pi}{9}}$  has been added to make it an element of SU(3). Further, the  $e_\beta$  coordinate of  $g_2 \cdot e_\alpha$  is obtained by evaluating



Since stacking the mirror image of a braid yields the identity, we know that we will obtain a unitary matrix. The braiding in the center of the diagram can be replaced by applying an *F*-move followed by an *R*-move. We obtain

$$\frac{\sqrt{\Delta_{\alpha}} \sqrt{\Delta_{\beta}}}{\theta(2,2,\alpha) \theta(2,2,\beta)} \sum_{i=0,2,4} \left\{ \begin{array}{c} 2 & 2 & i \\ 2 & 2 & 0 \end{array} \right\} \overline{\mathsf{R}}_{i}^{2,2} \left\{ \begin{array}{c} \beta \\ 2 & 2 & 0 \end{array} \right\}$$

It remains to compute the network



We do an *F*-move on the lower edge of the figure, labeled  $\alpha$ . It yields the first member of the equality below.

$$\sum_{j=0,2,4} \left\{ \begin{array}{c} 2 & 2 & j \\ 2 & 2 & \alpha \end{array} \right\} {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2 & \alpha \end{array} \right) {}_{2} \left( \begin{array}{c} \beta \\ 2 & 2$$

The figure to the right hand side of this equality is a tetrahedron. After replacing the 6j-symbol using Eq. (1) and simplifying, we then get

$$\frac{T \begin{bmatrix} 2 & 2 & i \\ 2 & 2 & \alpha \end{bmatrix} T \begin{bmatrix} 2 & 2 & i \\ 2 & 2 & \beta \end{bmatrix}}{\theta(2,2,i)}$$

We get a symmetric matrix whose  $(\alpha, \beta)$  coefficient is given by

$$\frac{\sqrt{\Delta_{\alpha}}\sqrt{\Delta_{\beta}}}{\theta(2,2,\alpha)\theta(2,2,\beta)} \sum_{i=0,2,4} \frac{\Delta_i \overline{R_i^{2,2}} T \begin{bmatrix} 2 & 2 & i \\ 2 & 2 & \alpha \end{bmatrix} T \begin{bmatrix} 2 & 2 & i \\ 2 & 2 & \beta \end{bmatrix}}{\theta(2,2,i)^2} ,$$

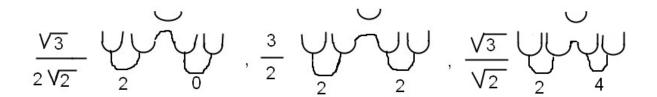
where we used that  $T\begin{bmatrix} 2 & 2 & i \\ 2 & 2 & 0 \end{bmatrix} = \theta(2, 2, i)$ . Up to the phase  $e^{\frac{i\pi}{9}}$  to make it a special unitary matrix, this matrix is the matrix  $G_2$  of Theorem 1. We obtain a unitary braid group representation  $B_4 \longrightarrow SU(3)$  whose image is  $\mathcal{G} = \langle G_1, G_2 \rangle$ .

In the next part, we show that the subgroup  $\mathcal{G}$  of SU(3) generated by the matrices  $G_1$  and  $G_2$  is isomorphic to the semidirect product  $\mathbb{Z}_9 \times \mathbb{Z}_3 \rtimes S_3$ , hence  $G_1$  and  $G_2$  generate D(9, 1, 1; 2, 1, 1).

## **3 Proof of the Theorem**

## 3.1 The abelian subgroup

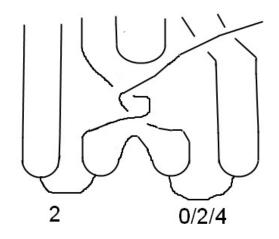
We introduce a new orthonormal basis with respect to the inner product defined by stacking a mirror diagram on top of another.



We consider the Hilbert space spanned over  $\mathbb{C}$  by these three vectors. In quantum computing words, the respective normalized diagrams are obtained by taking a quantum trit and an ancilla



by braiding anyons 4 and 5 and then doing a forced measurement to zero. It appears that doing a full twist in the center, that is braiding particules 3 and 6 twice leaves the Hilbert space invariant. After doing an additional two braidings between particules 6 and 7, and 7 and 8, the Hilbert space is still invariant.



After computing this action, we obtain the following unitary matrix

$$F = \begin{pmatrix} \frac{-1+i\sqrt{3}}{4} & \frac{\sqrt{2}}{4}(-1+i\sqrt{3}) & \frac{1-i\sqrt{3}}{4} \\ \frac{\sqrt{2}}{4}(-1+i\sqrt{3}) & 0 & \frac{\sqrt{2}}{4}(-1+i\sqrt{3}) \\ \frac{-1+i\sqrt{3}}{4} & \frac{\sqrt{2}}{4}(1-i\sqrt{3}) & \frac{1-i\sqrt{3}}{4} \end{pmatrix}$$

This matrix is  $G_1G_2 G_1^{-1} G_1^{-1}$ . Now consider

$$A = G_2 F G_2 F = (G_2 F)^2 = G_1 G_2^2 G_1^{-1},$$
  
$$B = G_1 G_2^{-2} G_1$$

and let  $\mathcal{N}$  be the subgroup of  $\mathcal{G}$  generated by A and B. The matrix A has order 9 and the matrix B has order 3. Since

$$B \neq A^3, B^2 \neq A^3, B \neq A^6, B^2 \neq A^6,$$

we know that

$$\langle A \rangle \cap \langle B \rangle = \{I\}$$

Moreover, we have

$$AB = BA$$

Then,

$$\mathcal{N}\simeq\mathbb{Z}_9 imes\mathbb{Z}_3$$

We now show that N is a normal subgroup of G. It suffices to show that for all the integers k and l, we have

$$G_1 A^k B^l G_1^{-1} \in \langle A, B \rangle$$
(2)

$$G_2 A^k B^l G_2^{-1} \in \langle A, B \rangle$$
(3)

Lemma 1. The following are equivalent.

$$\begin{aligned} (i) \ \forall \ l \ \in \mathbb{Z}, \ G_1 \ A^l \ G_1^{-1} \ \in < A, B > \\ (ii) \ \forall \ (k,l) \ \in \mathbb{Z}^2, \ G_1 \ A^l \ B^k \ G_1^{-1} \ \in < A, B > \end{aligned}$$

**Proof.** Suppose (*i*) holds. Write  $A^l B^k = (AB)^k A^{l-k}$ . Then,

$$G_{1} A^{l} B^{k} G_{1}^{-1} = G_{1} G_{1}^{2k} A^{l-k} G_{1}^{-1}$$
  
=  $G_{1}^{2k} G_{1} A^{l-k} G_{1}^{-1}$   
=  $(AB)^{k} G_{1} A^{l-k} G_{1}^{-1}$ 

Thus, in order to prove (2), it suffices to show that  $G_1 A G_1^{-1} \in \langle A, B \rangle$ . Since by definition,  $A = G_1 G_2^2 G_1^{-1}$  and since  $G_1^2 G_2^2 = G_2^2 G_1^2$ , we have

$$G_1 A G_1^{-1} = G_2^2$$

As  $G_2^2 = A^7 B^2$ , point (2) is proven. To have a normal subgroup, it remains to show (3). Since  $G_2^2 \in \langle A, B \rangle$ , we have

$$\begin{array}{lll} G_2 \, A^k \, B^l \, G_2^{-1} \, \in \, < A, B > & \Longleftrightarrow & G_2^2 \, G_2 \, A^k \, B^l \, G_2^{-1} \, \in \, < A, B > \\ & \Longleftrightarrow & G_2 \, A^{k-2} \, B^{l-1} \, G_2^{-1} \, \in \, < A, B > & (\clubsuit) \\ & \Leftrightarrow & G_2^{-2} \, G_2 \, A^k \, B^l \, G_2^{-1} \, \in \, < A, B > \\ & \Leftrightarrow & G_2 \, A^{k+2} \, B^{l+1} \, G_2^{-1} \, \in \, < A, B > & (\clubsuit) \end{array}$$

Equivalences ( $\clubsuit$ ) and ( $\bigstar$ ) show that  $G_2 A^k B^l G_2^{-1}$  belongs to  $\langle A, B \rangle$  for all  $(l,k) \in \mathbb{Z}^2$  is equivalent to  $G_2 A^t G_2^{-1}$  belongs to  $\langle A, B \rangle$  for every  $t \in \mathbb{Z}$ . Thus, it suffices to show that  $G_2 A G_2^{-1} \in \langle A, B \rangle$ . We have

$$G_2 A G_2^{-1} = G_2 G_1 G_2 G_2 G_1^{-1} G_2^{-1}$$
(4)

$$= G_1 G_2 G_1 G_2 G_1^{-1} G_2^{-1}$$
(5)

$$= G_1^2 G_2 G_1 G_1^{-1} G_2^{-1}$$
 (6)

$$= G_1^2 \tag{7}$$

$$= AB \tag{8}$$

Eq. (4) is by definition of *A* and Eq. (5) and (6) are obtained by using the braid relation. Hence, we are done with the proof of (3). So,  $\mathcal{N} \triangleleft \mathcal{G}$ .

#### 3.2 The symmetric group

We now describe the subgroup of  $\langle G_1, G_2 \rangle$  that is isomorphic to  $S_3$ . The product  $T_1 = G_1G_2G_1$  provides an element of order 2 in the group, which we take as one of the two generators of the symmetric group  $S_3$ . Since  $G_1$  has order 18, the product  $T_2 = G_2 G_1^9 G_2^{-1}$  provides another element of order 2 in the group. And so does  $G_2 G_1^2$ . Now set

$$T_3 = T_2 (G_2 G_1^2) T_2 = \begin{pmatrix} -1 & & \\ & -1 & \\ & & 1 \end{pmatrix}$$

and set

$$\mathcal{H} = < T_1, T_3 >$$

 $T_1$  and  $T_3$  generate a subgroup  $\mathcal{H}$  of  $\mathcal{G}$  isomorphic to  $S_3$ , whose two elements of order 3 are  $T_1 T_3$  and  $T_3 T_1$ . We checked using Mathematica that these two elements of order 3 are not one of

$$A^3, A^6, A^3 B, A^6 B, A^3 B^2, A^6 B^2, B, B^2$$

Then we must have

$$\mathcal{H} \cap \mathcal{N} = \{I\}$$

### 3.3 The semidirect product

We show that both  $G_1$  and  $G_2$  can be expressed as the product of an element of  $\mathcal{N}$  and an element of  $\mathcal{H}$ . This will suffice to show that  $\mathcal{G} \simeq \mathcal{N} \stackrel{\phi}{\rtimes} \mathcal{H}$ , where

$$\begin{aligned} \phi : & \mathcal{H} & \longrightarrow & Aut(\mathcal{N}) \\ & H & \longmapsto & \left( \begin{array}{cc} \mathcal{N} & \rightarrow & \mathcal{N} \\ & N & \mapsto & HNH^{-1} \end{array} \right) \end{aligned}$$

To show that  $G_1 \in \mathcal{N}.\mathcal{H}$ , it suffices to show that  $G_2^2 G_1 \in \mathcal{N}.\mathcal{H}$ . From  $(G_2 G_1 G_2)^2 = I$  and the braid relation, we derive  $(G_2 G_1)^3 = I$ . Then,

$$(G_2^2 G_1)^2 = G_2 G_2 G_1 G_2 G_2 G_1 = (G_2 G_1)^3 = I$$

So,  $G_2^2 G_1$  is an element of order 2 in  $\mathcal{G}$ . Moreover,  $G_2^2 G_1$  is not one of the three elements of order 2 of  $\mathcal{H}$ . Recall that in  $(\mathcal{N} \stackrel{\phi}{\rtimes} \mathcal{H}, \star)$ , we have

$$(N, H) \star (N, H) = (N H N H^{-1}, H^2)$$

So, if an element (N, H) has order 2, then H has order 2 in  $\mathcal{H}$  and N H N = H. Suppose we can write  $G_2^2 G_1 = N H$  with  $N \in \mathcal{N}$  and  $H \in \mathcal{H}$ . Impose the extra condition N H N = H. Then,

$$G_2^2 G_1 N = N H N = H$$

This implies  $N = (G_2^2 G_1)^{-1} H$ . Since  $\mathcal{N}$  does not contain any element of order 2, we see that H cannot be the identity matrix. Assume further that H has order 2. Then, H must be one of

$$T_1$$
 or  $T_3$  or  $T_3 T_1 T_3$ 

We found out that in the three cases,  $(G_2^2 G_1)^{-1} H$  has order 3. Out of the three possibilities, only

$$(G_2^2 G_1)^{-1} T_3$$

belongs to  $\mathcal{N}$  and is  $A^3 B$ . We get

$$G_2^2 G_1 = A^3 B T_3 \in \mathcal{N}.\mathcal{H}$$

And we conclude

$$G_1 = G_2^{-2} G_2^2 G_1 \in \mathcal{N}.\mathcal{H}$$

Similarly, we show that  $G_1^2 G_2 \in \mathcal{N}.\mathcal{H}$ . Again, since  $(G_1 G_2)^3 = I$ , this element has order 2. Also, it does not belong to  $\mathcal{H}$ . We then use the same strategy as before. If  $G_1^2 G_2 = NH$  with  $N \in \mathcal{N}$ ,  $H \in \mathcal{H}$  and NHN = H, then

$$N = (G_1^2 \, G_2)^{-1} \, H$$

With H of order 2, this time we find

$$N = (G_1^2 G_2)^{-1} T_3 T_1 T_3 = B^2$$

And so,

$$G_1^2 G_2 = B^2 T_3 T_1 T_3 \in \mathcal{N}.\mathcal{H}$$

Thus, since  $G_1^2 = AB$ , we get

$$G_2 = G_1^{-2} G_1^2 G_2 \in \mathcal{N}.\mathcal{H}$$

We are ready to conclude. Consider the map

$$\begin{array}{cccc} \mathcal{N} \stackrel{\phi}{\rtimes} \mathcal{H} & \stackrel{\psi}{\longrightarrow} & \mathcal{G} \\ (N,H) & \longmapsto & NH \end{array}$$

We have seen at the end of paragraph 3.2 that  $\mathcal{H} \cap \mathcal{N} = \{I\}$ . Hence  $\psi$  is an injective morphism of groups. By the discussion above, we have

$$G_1 = \psi(A^5 B^2, T_3)$$
  

$$G_2 = \psi(A^{-1} B, T_3 T_1 T_3)$$

It follows that

$$\langle G_1, G_2 \rangle \subseteq Im \psi \subseteq \langle G_1, G_2 \rangle$$

Hence the image of  $\psi$  is the whole group  $\mathcal{G}$  and

$$\mathcal{G}\simeq\mathcal{N}\stackrel{\phi}{
times}\mathcal{H}$$

We end this paper by giving a presentation for the group  $\mathcal{G}$ . We have

$$\mathcal{G} = \left\langle \begin{array}{ccc} A^9 = B^3 = T_1^2 = T_3^2 = (T_1 T_3)^3 = (T_3 T_1)^3 = I \\ A, B, T_1, T_3 & | & T_1 A T_1^{-1} = A, & T_3 A T_3^{-1} = A^7 B^2 \\ & T_1 B T_1^{-1} = A^6 B^2, & T_3 B T_3^{-1} = A^3 B^2 \end{array} \right\rangle$$

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## References

- L. Kauffmann and S. Lins, *Temperley-Lieb recoupling theory and invariants of 3-manifolds* Ann. Math. Studies, Vol 134, Princeton, NJ:Princeton Univ. Press 1994
- [2] P.O. Ludl, Systematic analysis of finite family symmetry groups and their applications to the lepton sector, Diploma thesis, University of Vienna 2009
- [3] P.O. Ludl, Comments on the classification of the finite subgroups of SU(3), J. Phys. A: Math. Theor. 44 255204 2011
- [4] G.A. Miller, H.F. Blichfeldt and L.E. Dickson, *Theory and Applications of Finite Groups*, John Wiley and Sons, New York 1916
- [5] Z. Wang, *Topological quantum computation*, CBMS monograph, Vol 112, American Mathematical Society 2010